

α -Logarithmic negativity

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Motivation: Entanglement and its manipulation

- Separable state: $\rho_{AB} = \sum_i p_i \rho_A^i \otimes \rho_B^i$.
- Entangled state: $\rho_{AB} \neq \sum_i p_i \rho_A^i \otimes \rho_B^i$.
- The most natural set of free operations for entanglement manipulation consists of **local operations and classical communication (LOCC)**, which has a complex structure [CLM⁺14].
- Entangled states cannot be created by LOCC.
- Inspired the resource theory framework: free states + free operations.
- The seminal ideas coming from it are influencing diverse areas: quantum thermodynamics, quantum coherence and superposition, non-Gaussianity, magic states...

Quantifying entanglement

- Entanglement is a key physical resource in quantum information, quantum computation, and quantum cryptography.
- A quantitative theory is highly desirable to fully exploit the power of entanglement.
- Entanglement measure E
 - **Faithfulness:** $E(\rho) = 0$ if and only if ρ is separable.
 - **LOCC monotonicity:** $E(\Lambda(\rho)) \leq E(\rho)$ for any $\Lambda \in \text{LOCC}$.
 - **Entanglement monotone, convexity, additivity**, etc.
- Zoo of ent. measures [PV07, Chr06].
- Information-processing task gives precise and operationally meaningful way to quantify a given physical resource.

- Due to the mathematical difficulty of dealing with separability and LOCC, it can be helpful to move beyond it.
- The **positive partial transpose** (PPT) criterion was proposed early [Per96, HHH96]

- A state ρ_{AB} is PPT if

$$T_B(\rho_{AB}) \geq 0,$$

where partial transpose map T_B for orthonormal basis $\{|i\rangle_B\}_i$ is defined as

$$T_B(Y_{AB}) \equiv \sum_{i,j} (I_A \otimes |i\rangle\langle j|_B) Y_{AB} (I_A \otimes |i\rangle\langle j|_B),$$

- If a state is separable, then it is PPT. Converse is not true in general.

Resource theory of NPT entanglement

- Rains proposed the **resource theory of non-positive partial transpose entanglement** [Rai99, Rai01]. That is, PPT states are free and non-PPT (NPT) states are resourceful.
- Free operations are the completely PPT preserving (C-PPT-P) channels. A bipartite channel $\mathcal{N}_{AB \rightarrow A'B'}$ is C-PPT-P if

$$T_{B'} \circ \mathcal{N}_{AB \rightarrow A'B'} \circ T_B \in \text{CP}$$

- Equivalently, the Choi operator $J_{AA'BB'}^{\mathcal{N}}$ for $\mathcal{N}_{AB \rightarrow A'B'}$ is PPT:

$$T_{BB'}(J_{AA'BB'}^{\mathcal{N}}) \geq 0$$

- Key result: All LOCC bipartite channels are C-PPT-P. So then we can use this relationship to find bounds for tasks in resource theory of entanglement by using resource theory of NPT entanglement.

Logarithmic negativity

- One of the most popular measures for resource theory of NPT entanglement is the **logarithmic negativity** [VW02, Ple05]:

$$E_N(\rho_{AB}) \equiv \log_2 \|T_B(\rho_{AB})\|_1,$$

- Entanglement monotone:

$$E_N(\rho_{AB}) \geq \sum_x p(x) E_N(\rho_{AB}^x)$$

where $\{p(x), \rho_{AB}^x\}_x$ is ensemble resulting from C-PPT-P instrument.

- Faithful on PPT states:

$$E_N(\rho_{AB}) = 0 \quad \iff \quad \rho_{AB} \in \text{PPT}$$

- It is neither convex nor monogamous.

- Recently, the κ -entanglement measure was defined as [WW18]

$$E_{\kappa}(\rho_{AB}) \equiv \log_2 \inf \{ \text{Tr}[S_{AB}] : \\ - S_{AB} \leq T_B(\rho_{AB}) \leq S_{AB}, T_B(S_{AB}) \geq 0 \}.$$

- Can be computed by semi-definite programming [WW18].
- Has a direct operational meaning in the resource theory of NPT entanglement, being equal to the exact PPT entanglement cost.
- It is an entanglement monotone, faithful on PPT states, but it is neither convex nor monogamous.
- Reduces to logarithmic negativity for all bipartite two-qubit states and bipartite bosonic Gaussian states.

Relation between log. negativity and κ -entanglement?

- Basic question is whether there is some relation between the two entanglement measures.
- **Result:** logarithmic negativity and κ -entanglement are extremes of a family of entanglement measures

Key quantities

- Consider

- ① $\alpha \geq 1$,
- ② a Hermitian operator $X \neq 0$, and
- ③ a positive semi-definite operator $\sigma \neq 0$:

- Then define

$$\mu_\alpha(X\|\sigma) \equiv \begin{cases} \left\| \sigma^{\frac{1-\alpha}{2\alpha}} X \sigma^{\frac{1-\alpha}{2\alpha}} \right\|_\alpha & \text{if } \text{supp}(X) \subseteq \text{supp}(\sigma) \\ +\infty & \text{else} \end{cases},$$

$$\nu_\alpha(X\|\sigma) \equiv \log_2 \mu_\alpha(X\|\sigma),$$

- Key properties of $\mu_\alpha(X\|\sigma)$ and $\nu_\alpha(X\|\sigma)$ can be derived from [Bei13] and [Hia16]

Definition (α -logarithmic negativity)

Let ρ_{AB} be a bipartite state. Its α -**logarithmic negativity** defined as

$$E_N^\alpha(\rho_{AB}) \equiv \inf_{\sigma_{AB} \in \text{PPT}(A:B)} \nu_\alpha(T_B(\rho_{AB}) \parallel \sigma_{AB}),$$

where $\text{PPT}(A : B)$ is the set of PPT states:

$$\text{PPT}(A : B) \equiv \{\sigma_{AB} : \sigma_{AB}, T_B(\sigma_{AB}) \geq 0, \text{Tr}[\sigma_{AB}] = 1\}.$$

We now discuss properties of the α -logarithmic negativities.

Proposition

Let ρ_{AB} be a bipartite quantum state, and let $1 \leq \alpha \leq \beta$. Then

$$E_N(\rho_{AB}) \leq E_N^\alpha(\rho_{AB}) \leq E_N^\beta(\rho_{AB}).$$

Follows from a simple generalization of [Bei13, Theorem 7]

Lemma

Let $X \neq 0$ be a Hermitian operator, and let σ be a positive definite operator. Then the following inequality holds for all $\beta > \alpha > 1$:

$$\frac{\alpha}{\alpha - 1} [\nu_\alpha(X\|\sigma) - \log_2 \|X\|_1] \leq \frac{\beta}{\beta - 1} [\nu_\beta(X\|\sigma) - \log_2 \|X\|_1]$$

which implies

$$\nu_\alpha(X\|\sigma) \leq \nu_\beta(X\|\sigma).$$

Proposition

Let ρ_{AB} be a bipartite quantum state. Then

$$\lim_{\alpha \rightarrow 1} E_N^\alpha(\rho_{AB}) = E_N(\rho_{AB}).$$

Definition (Max-logarithmic negativity)

For bipartite state ρ_{AB} , max-logarithmic negativity $E_N^{\max}(\rho_{AB})$ defined as

$$E_N^{\max}(\rho_{AB}) \equiv \inf_{\sigma_{AB} \in \text{PPT}(A:B)} \nu_{\infty}(T_B(\rho_{AB}) \parallel \sigma_{AB}),$$

and

$$\nu_{\infty}(X \parallel \sigma) = D_{\max}(X \parallel \sigma) = \log_2 \left\| \sigma^{-1/2} X \sigma^{-1/2} \right\|_{\infty}.$$

Proposition

Let ρ_{AB} be a bipartite quantum state. Then

$$E_{\kappa}(\rho_{AB}) = E_N^{\max}(\rho_{AB}) = \lim_{\alpha \rightarrow \infty} E_N^{\alpha}(\rho_{AB}).$$

Proposition

If ρ_{AB} satisfies the condition $T_B(|T_B(\rho_{AB})|) \geq 0$, then all α -logarithmic negativities are equal; i.e., the following equality holds for all $\alpha \geq 1$:

$$E_N^\alpha(\rho_{AB}) = E_N(\rho_{AB}).$$

All of the following satisfy $T_B(|T_B(\rho_{AB})|) \geq 0$:

- 1 pure states [ADMVW02],
- 2 two-qubit states [Ish04],
- 3 Werner states [APE03], and
- 4 bosonic Gaussian states [APE03]

Thus, conclude that collapse above holds for such states.

Definition

A **C-PPT-P quantum instrument** consists of the collection

$$\{\mathcal{N}_{AB \rightarrow A'B'}^x\}_x,$$

where

- 1 each $\mathcal{N}_{AB \rightarrow A'B'}^x$ is CP,
- 2 the map $T_{B'} \circ \mathcal{N}_{AB \rightarrow A'B'}^x \circ T_B$ is CP, and
- 3 the sum map $\sum_x \mathcal{N}_{AB \rightarrow A'B'}^x$ is TP.

Theorem (Entanglement monotone)

Let $\{\mathcal{N}_{AB \rightarrow A'B'}^x\}_x$ be a C-PPT-P quantum instrument, and ρ_{AB} a bipartite state. Then α -logarithmic negativity is an **entanglement monotone**; i.e., the following inequality holds for all $\alpha \geq 1$:

$$E_N^\alpha(\rho_{AB}) \geq \sum_{x:p(x)>0} p(x) E_N^\alpha(\rho_{A'B'}^x),$$

where

$$p(x) \equiv \text{Tr}[\mathcal{N}_{AB \rightarrow A'B'}^x(\rho_{AB})],$$
$$\rho_{A'B'}^x \equiv \frac{1}{p(x)} \mathcal{N}_{AB \rightarrow A'B'}^x(\rho_{AB}).$$

Entanglement monotone proof

Follows from a simple generalization of [Bei13, Theorem 6]

Lemma

Let $X \neq 0$ be a Hermitian operator, and let σ be a positive definite operator. Let \mathcal{P} be a positive and trace-non-increasing map. Then the following inequality holds for all $\alpha \geq 1$:

$$\nu_\alpha(X \parallel \sigma) \geq \nu_\alpha(\mathcal{P}(X) \parallel \mathcal{P}(\sigma)).$$

Entanglement monotone proof

Entanglement monotone property also follows from

Lemma

$$\text{Let } Y_{XB} \equiv \sum_x p(x) |x\rangle\langle x|_X \otimes Y_B^x, \quad \sigma_{XB} \equiv \sum_x q(x) |x\rangle\langle x|_X \otimes \sigma_B^x,$$

where $\{Y_B^x\}_x$ is a set of Hermitian operators, $\{p(x)\}_x$ is a probability distribution, $\{\sigma_B^x\}$ is a set of positive definite operators, and $\{q(x)\}_x$ is a set of strictly positive reals. Then for $\alpha \geq 1$,

$$\nu_\alpha(Y_{XB} \| \sigma_{XB}) \geq \sum_x p(x) \nu_\alpha(Y_B^x \| \sigma_B^x) + \left(\frac{\alpha - 1}{\alpha} \right) D(p \| q),$$

where $D(p \| q) := \sum_x p(x) \log_2(p(x)/q(x))$ is the classical relative entropy.

Proposition

Let ρ_{AB} be a bipartite state and $\alpha \geq 1$. Then the α -logarithmic negativity $E_N^\alpha(\rho_{AB})$ can be calculated by convex optimization.

The above proposition follows from

Lemma

Let $X \neq 0$ be a Hermitian operator, and let σ be a positive definite operator. Then for all $\alpha \geq 1$, the following function is convex:

$$\sigma \mapsto [\mu_\alpha(X \parallel \sigma)]^\alpha.$$

The above lemma follows directly from [Hia16, Theorem 5.2] (see also brief remarks stated before [Hia16, Theorem 5.3])

Proposition (Faithfulness)

Let ρ_{AB} be a bipartite quantum state, and let $\alpha \geq 1$. Then

$$E_N^\alpha(\rho_{AB}) \geq 0$$

and

$$E_N^\alpha(\rho_{AB}) = 0 \quad \iff \quad \rho_{AB} \in \text{PPT}(A : B).$$

No convexity

By picking

$$\begin{aligned}\rho_{AB}^1 &\equiv \Phi_{AB}^2, & \rho_{AB}^2 &\equiv \frac{1}{2} (|00\rangle\langle 00|_{AB} + |11\rangle\langle 11|_{AB}), \\ \bar{\rho}_{AB} &\equiv \frac{1}{2} (\rho_{AB}^1 + \rho_{AB}^2),\end{aligned}$$

we find for all $\alpha \in [1, \infty]$ that

$$\begin{aligned}E_N^\alpha(\rho_{AB}^1) &= 1, & E_N^\alpha(\rho_{AB}^2) &= 0, \\ E_N^\alpha(\bar{\rho}_{AB}) &= \log_2 \frac{3}{2},\end{aligned}$$

which implies no convexity:

$$E_N^\alpha(\bar{\rho}_{AB}) > \frac{1}{2} [E_N^\alpha(\rho_{AB}^1) + E_N^\alpha(\rho_{AB}^2)].$$

Definition of monogamy

An entanglement measure E is monogamous [CKW00, Ter04, KW04] if the following inequality holds for all tripartite states ρ_{ABC} :

$$E(\rho_{A:B}) + E(\rho_{A:C}) \leq E(\rho_{A:BC}),$$

where the bipartition is indicated by a colon.

- As a consequence of the counterexample given in [WW18, Proposition 7], it follows that the α -logarithmic negativity is not generally monogamous for any choice of $\alpha \in [1, \infty]$.
- Indeed, consider the following state of three qubits:

$$|\psi\rangle_{ABC} \equiv \frac{1}{2} \left(|000\rangle_{ABC} + |011\rangle_{ABC} + \sqrt{2}|110\rangle_{ABC} \right).$$

Then

$$E_N^\alpha(\psi_{A:B}) + E_N^\alpha(\psi_{A:C}) > E_N^\alpha(\psi_{A:BC}),$$

Additivity?

Given tensor-product state $\omega_{A_1A_2:B_1B_2} = \rho_{A_1B_1} \otimes \tau_{A_2B_2}$.

- Logarithmic negativity is **additive** [VW02]:

$$E_N(\omega_{A_1A_2:B_1B_2}) = E_N(\rho_{A_1:B_1}) + E_N(\tau_{A_2:B_2}).$$

- So is κ -entanglement [WW18]:

$$E_\kappa(\omega_{A_1A_2:B_1B_2}) = E_\kappa(\rho_{A_1:B_1}) + E_\kappa(\tau_{A_2:B_2}).$$

- For α -logarithmic negativity, subadditivity holds for $\alpha \in (1, \infty)$:

$$E_N^\alpha(\omega_{A_1A_2:B_1B_2}) \leq E_N^\alpha(\rho_{A_1:B_1}) + E_N^\alpha(\tau_{A_2:B_2}).$$

What about superadditivity?

Generalization to channels

- **Logarithmic negativity of a channel** defined as [HW01]

$$E_N(\mathcal{N}) \equiv \log_2 \|T_B \circ \mathcal{N}_{A \rightarrow B}\|_{\diamond}.$$

- **Diamond norm** of Herm.-preserving map $\mathcal{M}_{A \rightarrow B}$ defined as [Kit97]

$$\|\mathcal{M}_{A \rightarrow B}\|_{\diamond} \equiv \sup_{\psi_{RA}} \|\mathcal{M}_{A \rightarrow B}(\psi_{RA})\|_1,$$

- Can write the logarithmic negativity of a quantum channel as an optimized version of the logarithmic negativity of quantum states:

$$E_N(\mathcal{N}) = \sup_{\psi_{RA}} E_N(\omega_{RB}),$$

where $\omega_{RB} \equiv \mathcal{N}_{A \rightarrow B}(\psi_{RA})$.

Definition (α -log. negativity of a channel)

The α -logarithmic negativity of a quantum channel is defined for $\alpha \geq 1$ as

$$E_N^\alpha(\mathcal{N}) = \sup_{\psi_{RA}} E_N^\alpha(\omega_{RB}),$$

with $\omega_{RB} \equiv \mathcal{N}_{A \rightarrow B}(\psi_{RA})$.

Recover **κ -entanglement of a channel** as a special case ($\alpha \rightarrow \infty$), which is equal to the exact PPT simulation cost of a quantum channel (parallel or sequential simulation) [WW18]

Conclusion

- α -logarithmic negativities are a family of entanglement monotones that include logarithmic negativity and κ -entanglement
- Entanglement monotone property follows from techniques of [Bei13]
- They are ordered, faithful, and computable by convex optimization
- They are neither convex nor monogamous

- The concept put forward here can be generalized.
- Idea is to compare an unphysical object versus a physical one to obtain a useful information, distinguishability, or entanglement measure
- In our case, we compare the partial transpose of ρ_{AB} to the set of PPT states using $\nu_\alpha(X\|\sigma)$.

Going forward

Yesterday, we saw an interesting re-expression of diamond norm, using the same concept:

The slide is titled "Asymptotic Continuity" in a purple header. It contains the following text and equations:

$D_{\mathfrak{F}}$ and $E_{\mathfrak{F}}$ are asymptotically continuous.

$$|D_{\mathfrak{F}}(\mathcal{N}_{A \rightarrow B}) - D_{\mathfrak{F}}(\mathcal{M}_{A \rightarrow B})| \leq f(\|\mathcal{N} - \mathcal{M}\|_{\diamond}) \log |AB|$$
$$\lim_{\epsilon \rightarrow 0^+} f(\epsilon) = 0$$
$$\|\mathcal{N} - \mathcal{M}\|_{\diamond} := \max_{\psi_{RA}} \|\mathcal{N}_{A \rightarrow B}(\psi_{RA}) - \mathcal{M}_{A \rightarrow B}(\psi_{RA})\|_1$$

Key observation:

$$\log_2 \|\mathcal{N} - \mathcal{M}\|_{\diamond} - 1 = \min_{\mathcal{E} \in \text{CPTP}(A \rightarrow B)} D_{\max}(\mathcal{N} - \mathcal{M} \| \mathcal{E})$$

The last equation is enclosed in a red rectangular box.

Thanks!



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