

An endpoint Strichartz estimate in spherical coordinates

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Abstract

We study Strichartz estimates in spherical coordinates for dispersive equations which are defined by spherically symmetric pseudo-differential operators. We extend the recent results in [7, 11] to include more general class of dispersive equations. We use a bootstrapping argument based on various weighted Strichartz estimates.

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1 Introduction

In this paper we consider the Cauchy problem of linear dispersive equations:

$$iu_t - \omega(|\nabla|)u = 0 \text{ in } \mathbb{R}^{1+n}, \quad u(0) = \varphi \text{ in } \mathbb{R}^n, \quad n \geq 2 \quad (1.1)$$

where $|\nabla| = \sqrt{-\Delta}$ and the operator $\omega(|\nabla|)$ is the pseudo-differential operator of which multiplier is $\omega(|\xi|)$. We will work with $\omega \in C[0, \infty) \cap C^\infty(0, \infty)$ which satisfies the following properties:

- (i) $\omega'(\rho) > 0$, and either $\omega''(\rho) > 0$ or $\omega''(\rho) < 0$ for all $\rho > 0$,
- (ii) $\omega^{(k)}(\rho_1) \sim \omega^{(k)}(\rho_2)$, $k = 1, 2$ for $0 < \rho_1 < \rho_2 < 2\rho_1$,
- (iii) $\rho|\omega^{(k+1)}(\rho)| \lesssim |\omega^{(k)}(\rho)|$ for all $k \geq 1$ and $\rho > 0$.

Typical examples of ω are ρ^a ($0 < a \neq 1$), $\sqrt{1 + \rho^2}$, $\rho\sqrt{1 + \rho^2}$, and $\frac{\rho}{\sqrt{1 + \rho^2}}$ which describe the Schrödinger type equations (see [12] for $a < 2$), Klein-Gordon or semirelativistic [8], iBq, and imBq equations. (For the last two see [3] and references therein.)

The solution can formally be written by

$$u(t, x) = \frac{1}{(2\pi)^n} \int_{\mathbb{R}^n} e^{i(x \cdot \xi - t\omega(|\xi|))} \widehat{\varphi}(\xi) d\xi.$$

Here $\widehat{\varphi}$ is the Fourier transform of φ defined by $\int_{\mathbb{R}^n} e^{-ix \cdot \xi} \varphi(x) dx$. In [6] the standard Strichartz estimate in $L_t^q L_x^p$ was considered with ω satisfying (i), (ii), (iii) and the following was shown: if $n \geq 1$ and the pair (q, p) satisfies that $2 \leq q, p \leq \infty$, $\frac{1}{q} \leq \frac{n}{2}(\frac{1}{2} - \frac{1}{p})$ and $(q, p) \neq (2, \infty)$, then

$$\| |\nabla|^s \mathcal{D}_\omega^{s_1, s_2} u \|_{L_t^q L_x^p} \lesssim \|\varphi\|_{L_x^2} \quad (1.2)$$

for $s_1 = \frac{1}{q} - s_2$, $s_2 = \frac{1}{nq}$ and $s = \frac{2}{q} - n(\frac{1}{2} - \frac{1}{p})$, where $\mathcal{D}_\omega^{s_1, s_2}$ is a pseudo-differential operator whose symbol is

$$\left(\frac{\omega'(|\xi|)}{|\xi|} \right)^{s_1} |\omega''(|\xi|)|^{s_2}.$$

In this paper we study the estimate (1.2) by making use of mixed norm spaces given in the spherical coordinates. For this purpose we use the time-space norm given by

$$\|f\|_{L_r^p L_\sigma^\ell} = \left(\int_0^\infty \left(\int_{S^{n-1}} |f(r\sigma)|^\ell d\sigma \right)^{\frac{p}{\ell}} r^{n-1} dr \right)^{\frac{1}{p}}, \quad 1 \leq p, \ell \leq \infty.$$

For simplicity we denoted the spaces $L^p(r^{n-1}dr)$ by L_r^p . Clearly $\|f\|_{L_x^p} = \|f\|_{L_r^p L_\sigma^p}$. Then let us define several function spaces of Sobolev type. Let Δ_σ be the Laplace-Beltrami operator defined on the unit sphere and set $D_\sigma = \sqrt{1 - \Delta_\sigma}$. For $|s| < n/p$, $\gamma \in \mathbb{R}$, we denote by $\dot{H}_r^{s,p} H_\sigma^{\gamma,\ell}$ the space

$$\left\{ f \in \mathcal{S}' : \|f\|_{\dot{H}_r^{s,p} H_\sigma^{\gamma,\ell}} \equiv \|\nabla|^s D_\sigma^\gamma f\|_{L_r^p L_\sigma^\ell} < \infty \right\}.$$

It should be noted that C_0^∞ is dense in $\dot{H}_r^{s,p} H_\sigma^{\gamma,\ell}$ since $|s| < n/p$ (see [2]). We also use spaces equipped with the time-space norm

$$\|v\|_{L_t^q \dot{H}_r^{s,p} H_\sigma^{\gamma,\ell}} = \left(\int_{\mathbb{R}} \|v(\cdot, t)\|_{\dot{H}_r^{s,p} H_\sigma^{\gamma,\ell}}^q dt \right)^{\frac{1}{q}}, \quad 1 \leq q \leq \infty.$$

If $\ell = 2$, we use a simplified notation $\dot{H}_r^s H_\sigma^\gamma$ for $\dot{H}_r^{s,2} H_\sigma^{\gamma,2}$.

It is well known that the range of p, q for (1.2) can not be extended as it can be shown by Knapp's example. There have been results [7, 15, 16] which extend the range by allowing loss of angular regularity (also see [9, 13, 14] for related results). Such results have been proven to be useful in the study of various nonlinear equations [1, 7]. Recently in [11] the authors showed that if $n \geq 2$, $\frac{1}{q} < (n-1)(\frac{1}{2} - \frac{1}{p})$ or $(q, p) = (\infty, 2)$ for $p, q \geq 2$ and $(q, p) \neq (\infty, \infty), (2, \infty)$, then

$$\|\nabla|^s u\|_{L_t^q L_x^p} \lesssim \|\varphi\|_{L_r^2 H_\sigma^\gamma} \quad (1.3)$$

for $\omega(\rho) = \rho^a, a > 0$, $s = \frac{a}{q} - n(\frac{1}{2} - \frac{1}{p})$ and $\gamma \geq 1/q$. They utilized Rodnianski's argument in [15] and weighted Strichartz estimates (see [5, 7, 1]).

In this short note we show that the estimate (1.3) can be extended to include more general ω and the angular regularity can be improved (see

Proposition 3.1 below). For simplicity we consider only the endpoint case $q = 2$ since the full estimate can be obtained by interpolation with the trivial estimate $\|u\|_{L_t^\infty L_x^2} \lesssim \|\varphi\|_{L_x^2}$ or the estimates in Theorem 1.7 of [11]. The novelty here is the use of bootstrapping to extend the range of (1.4).

The following is our main result.

Theorem 1.1. *Suppose that ω satisfies the conditions (i)–(iii). Let $n \geq 3$, $\frac{2(n-1)}{n-2} < p < \frac{2n}{n-2}$ and $s_0 = \frac{1}{2} - \frac{n-1}{(n-2)p}$. Then for sufficiently small $\varepsilon > 0$ we have*

$$\| |\nabla|^s \mathcal{D}_\omega^{s_1, s_2} u \|_{L_t^2 L_x^p} \lesssim \|\varphi\|_{L_r^2 H^\gamma} \quad (1.4)$$

for $s = \frac{n}{p} - \frac{n-2}{2}$, $s_1 = \frac{1}{2} - s_0 - \varepsilon$, $s_2 = s_0 + \varepsilon$ and $\gamma > \frac{1}{2} - ns_0$.

If $\omega(\rho) = \rho^a$, $0 < a \neq 1$, then since $\omega'(\rho)/\rho \sim \omega''(\rho) \sim \rho^{a-2}$, we get (1.3) with γ as in Theorem 1.1. We will not pursue the optimality of angular regularity, which is another interesting issue.

For the proof of the theorem we use bootstrapping argument based on the Sobolev inequality and weighted Strichartz estimates in spherical coordinates. We will start bootstrapping from the endpoint Strichartz estimate. Once we have an endpoint estimate (1.4) for $p \neq \infty^1$, making use of Sobolev inequalities (2.1), (2.2) and (2.3), we get the estimate (1.4) for $p = p_k$, $k = 1, 2, 3, \dots$, successively. The sequence p_k decreasingly converges to $\frac{2(n-1)}{n-2}$. Regardless of ω , by this argument we can get estimate (1.4) arbitrarily close to $p = \frac{2(n-1)}{n-2}$ and thus via interpolation we also get the estimate $\| |\nabla|^s \mathcal{D}_\omega^{s_1, s_2} u \|_{L_t^q L_x^p} \lesssim \|\varphi\|_{L_r^2 H^\gamma}$ for $p, q \geq 2$ satisfying $\frac{1}{q} < (n-1)(\frac{1}{2} - \frac{1}{p})$. There is an ε -loss involved in s_1, s_2 which results from interpolations of the estimates for $p < \frac{2(n-1)}{n-2}$. But there is no loss if we impose additional condition that $\omega'(\rho) \sim \rho|\omega''(\rho)|$. If one can obtain an estimate on the critical line $L = \{(1/p, 1/q) : \frac{1}{q} = (n-1)(\frac{1}{2} - \frac{1}{p})\}$, such loss can be removed.

Finally, we make a remark on the case of the wave equation in which $\omega(\rho) = \rho$. Using the known endpoint Strichartz estimate for $n \geq 4$ ([10])

¹This is why we work with $n \geq 3$. Actually if $n = 2$, then Strichartz estimate in spherical coordinates is worse than the standard one.

for the wave equation, one may apply the bootstrapping argument of this note to get the same result as of Sterbenz [15]. However it is not possible to remove the ε -loss in the angular regularity this way. For this it seems that one needs to obtain estimates along the critical line L .

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2 Weighted estimates

We first recall Sobolev inequality which was introduced in [4] and extended in [7]. Let $0 < a < \frac{n-1}{2}$ and $\alpha \leq \frac{n-1}{2} - a$. Then

$$\| |x|^a |\nabla|^{a-\frac{n}{2}} D_\sigma^\alpha f \|_{L_r^\infty L_\sigma^2} \lesssim \|f\|_{L_x^2}, \quad (2.1)$$

In [2], the cases $f \in L_x^p$, $1 \leq p < 2$ were treated. Using complex interpolation between (2.1) and the trivial estimates $\|f\|_{L_r^\infty L_\sigma^\infty} \lesssim \|f\|_{L_r^\infty L_\sigma^\infty}$ and $\|f\|_{L_r^\infty L_\sigma^2} \lesssim \|f\|_{L_r^\infty L_\sigma^2}$, we get for $2 \leq \ell \leq \infty$

$$\| |x|^{\frac{2a}{\ell}} |\nabla|^{\frac{2}{\ell}(a-\frac{n}{2})} D_\sigma^{\frac{2\alpha}{\ell}} f \|_{L_r^\infty L_\sigma^\ell} \leq C_0 \|f\|_{L_x^\ell}, \quad (2.2)$$

$$\| |x|^{\frac{2a}{\ell}} |\nabla|^{\frac{2}{\ell}(a-\frac{n}{2})} D_\sigma^{\frac{2\alpha}{\ell}} f \|_{L_r^\infty L_\sigma^2} \leq C_0 \|f\|_{L_r^\ell L_\sigma^2}. \quad (2.3)$$

Replacing f of (2.2) with u and applying endpoint Strichartz estimate (1.2), we obtain the following.

Lemma 2.1. *Let $n \geq 3$, $0 < a < \frac{n-1}{2}$ and $\alpha \leq \frac{n-1}{2} - a$. Then*

$$\| |x|^{a(1-\frac{2}{n})} |\nabla|^{(a-\frac{n}{2})(1-\frac{2}{n})} \mathcal{D}_\omega^{\frac{n-1}{2n}, \frac{1}{2n}} D_\sigma^{\alpha(1-\frac{2}{n})} u \|_{L_t^2 L_r^\infty L_\sigma^{\frac{2n}{n-2}}} \leq C'_0 \|\varphi\|_{L_x^2}. \quad (2.4)$$

Proof. From (2.2) with $\ell = \frac{2n}{n-2}$ and Hölder's inequality it follows that

$$\begin{aligned} & \| |x|^{a(1-\frac{2}{n})} |\nabla|^{(a-\frac{n}{2})(1-\frac{2}{n})} \mathcal{D}_\omega^{\frac{n-1}{2n}, \frac{1}{2n}} D_\sigma^{\alpha(1-\frac{2}{n})} u \|_{L_t^2 L_r^\infty L_\sigma^{\frac{2n}{n-2}}} \\ & \leq C_0 \| \mathcal{D}_\omega^{\frac{n-1}{2n}, \frac{1}{2n}} u \|_{L_t^2 L_r^{\frac{2n}{n-2}} (L_\sigma^2 \cap L_\sigma^{\frac{2n}{n-2}})} \leq c_n C_0 \| \mathcal{D}_\omega^{\frac{n-1}{2n}, \frac{1}{2n}} u \|_{L_t^2 L_x^{\frac{2n}{n-2}}}, \end{aligned}$$

where c_n depends only on n . Then the estimate (1.2) with $\ell = \frac{2n}{n-2}$ gives (2.4). \square

We will use the following $L_t^2 L_x^2$ estimate.

Lemma 2.2. *Let $-n/2 < b < -1/2$ and $\beta \leq -\frac{1}{2} - b$. Then we have*

$$\| |x|^b |\nabla|^{b+1} \mathcal{D}_\omega^{\frac{1}{2},0} D_\sigma^\beta u \|_{L_t^2 L_x^2} \leq C_1 \|\varphi\|_{L_x^2}. \quad (2.5)$$

For (2.5) we refer to [1] and also to [5, 7] for earlier versions.

3 Proof of Theorem 1.1

In this section we prove Theorem 1.1 by showing Proposition 3.1. It will be shown via bootstrapping argument which makes use of weighted Strichartz estimates introduced in the previous section.

Proposition 3.1. *Let ω satisfy (i) – (iii) and $n \geq 3$. Then, for $\frac{2(n-1)}{n-2} < p < \infty$,*

$$\| |\nabla|^{-\frac{n-2}{2} + \frac{n}{p}} \mathcal{D}_\omega^{\frac{1}{2},0} D_\sigma^{\frac{n-2}{2} - \frac{n-1}{p}} u \|_{L_t^2 L_r^p L_\sigma^2} \lesssim \|\varphi\|_{L_x^2}. \quad (3.1)$$

Assuming this for the moment, we prove Theorem 1.1. In fact, using (3.1) and Sobolev embedding $(H_\sigma^{\ell, (n-1)(1/2-1/p)} \hookrightarrow L_\sigma^p)$ on the unit sphere, we have for $\frac{2(n-1)}{n-2} < p < \infty$

$$\| |\nabla|^s \mathcal{D}_\omega^{\frac{1}{2},0} u \|_{L_t^2 L_x^p} \lesssim \|\varphi\|_{L_r^2 H_\sigma^{\frac{1}{2}}},$$

where $s = \frac{n}{p} - \frac{n-2}{2}$. Now, from the endpoint estimate of (1.2) we have

$$\| |\nabla|^s \mathcal{D}_\omega^{\frac{1}{2} - \frac{1}{2n}, \frac{1}{2n}} u \|_{L_t^2 L_x^{\frac{2n}{n-2}}} \lesssim \|\varphi\|_{L_r^2 H_\sigma^0}.$$

Interpolation between these two estimates gives the desired result.

Now it remains to prove Proposition 3.1 .

3.1 Bootstrapping

We start with interpolating (2.4) with $\alpha = (n-1)/2 - a$ and (2.5) with $\beta = -1/2 - b$. So, we have for $2 \leq p_1 \leq \infty$

$$\| |x|^{c_1} |\nabla|^{d_1} \mathcal{D}_\omega^{\delta_1, \delta'_1} D_\sigma^{\gamma_1} u \|_{L_t^2 L_r^{p_1} (L_\sigma^2 \cap L_\sigma^{\ell_1})} \leq \mathcal{C}_1 \|\varphi\|_{L_x^2}, \quad (3.2)$$

where $\mathcal{C}_1 = C_0'^{1-\theta_1} C_1^{\theta_1}$, $\theta_1 = 2/p_1$, and

$$\begin{aligned} c_1 &= a(1 - 2/n)(1 - 2/p_1) + 2b/p_1, \\ d_1 &= c_1 - (n-2)/2 + n/p_1, \\ \delta_1 &= \frac{1}{n} \left(\frac{n-1}{2} + \frac{1}{p_1} \right), \quad \delta'_1 = \frac{1}{n} \left(\frac{1}{2} - \frac{1}{p_1} \right), \\ \gamma_1 &= \frac{n^2 - 3n + 2}{2n} - \frac{n^2 - 2n + 2}{np_1} - c_1, \\ \frac{1}{\ell_1} &= \frac{n-2}{2n} + \frac{1}{np_1}. \end{aligned}$$

We call it the first stage estimate and will proceed similarly by combining the resulting estimate and (2.5). At every stage we choose the indices $c_k, k \geq 1$ to be 0. In view of the range of a and b we can take $c_1 = 0$ when p_1 satisfies that $2 + \frac{2n}{(n-1)(n-2)} = p_0 < p_1 < \infty$. In particular, such p_1 gives

$$\| |\nabla|^{d_1} \mathcal{D}_\omega^{\delta_1, \delta'_1} D_\sigma^{\gamma_1} u \|_{L_t^2 L_r^{p_1} (L_\sigma^2 \cap L_\sigma^{\ell_1})} \leq \mathcal{C}_1 \|\varphi\|_{L_x^2}. \quad (3.3)$$

We use it in the place of Strichartz estimate (1.2). Now, from the estimates (2.3) and (3.3) it follows that

$$\begin{aligned} & \| |x|^{\frac{2a}{p_1}} |\nabla|^{\frac{2}{p_1}(a - \frac{n}{2}) + d_1} \mathcal{D}_\omega^{\delta_1, \delta'_1} D_\sigma^{\frac{2\alpha}{p_1} + \gamma_1} u \|_{L_t^2 L_r^\infty L_\sigma^2} \\ & \leq C_0 \| |\nabla|^{d_1} \mathcal{D}_\omega^{\delta_1, \delta'_1} D_\sigma^{\gamma_1} u \|_{L_t^2 L_r^{p_1} L_\sigma^2} \\ & \leq C_0 \mathcal{C}_1 \|\varphi\|_{L_x^2}, \end{aligned}$$

where a and α are given in Lemma 2.1. Interpolating this with (2.5) for $\alpha = \frac{n-1}{2} - a$ and $\beta = -b - \frac{1}{2}$, we have the second stage estimate: for $2 \leq p_2 \leq \infty$

$$\| |x|^{c_2} |\nabla|^{d_2} \mathcal{D}_\omega^{\delta_2, \delta'_2} D_\sigma^{\gamma_2} u \|_{L_t^2 L_r^{p_2} L_\sigma^2} \leq \mathcal{C}_2 \|\varphi\|_{L_x^2},$$

where $\mathcal{C}_2 = (C_0 \mathcal{C}_1)^{1-\theta_2} C_1^{\theta_2}$, $\theta_2 = 2/p_2$ and

$$\begin{aligned} c_2 &= \frac{2a}{p_1} \left(1 - \frac{2}{p_2}\right) + \frac{2b}{p_2} < \frac{2a}{p_0} \left(1 - \frac{2}{p_2}\right) + \frac{2b}{p_2}, \\ d_2 &= c_2 + \left(-\frac{n}{p_1} + d_1\right) \left(1 - \frac{2}{p_2}\right) + \frac{2}{p_2}, \\ \delta_2 &= \delta_1 \left(1 - \frac{2}{p_2}\right) + \frac{1}{p_2}, \quad \delta'_2 = \delta'_1 \left(1 - \frac{2}{p_2}\right), \\ \gamma_2 &= \left(\gamma_1 + \frac{n-1}{p_1}\right) \left(1 - \frac{2}{p_2}\right) - \frac{1}{p_2} - c_2. \end{aligned}$$

If $2 + \frac{p_0}{n-1} < p_2 < \infty$, then by suitable choices of a, b , we can make $c_2 = 0$.

Thus $d_2 = -\frac{n-2}{2} + \frac{n}{p_2}$, $\gamma_2 > 0$ and we also have

$$\| |\nabla|^{d_2} \mathcal{D}_\omega^{\delta_2, \delta'_2} D_\sigma^{\gamma_2} u \|_{L_t^2 L_r^{p_2} L_\sigma^2} \leq \mathcal{C}_2 \|\varphi\|_{L_x^2}.$$

Repeating this procedure k times, we obtain

$$\| |x|^{c_k} |\nabla|^{d_k} \mathcal{D}_\omega^{\delta_k, \delta'_k} D_\sigma^{\gamma_k} u \|_{L_t^2 L_r^{p_k} L_\sigma^2} \leq \mathcal{C}_k \|\varphi\|_{L_x^2},$$

where $\mathcal{C}_k = (C_0 \mathcal{C}_{k-1})^{1-\theta_k} C_1^{\theta_k}$, $\theta_k = 2/p_k$ and

$$\begin{aligned} c_k &= \frac{2a}{p_{k-1}} \left(1 - \frac{2}{p_k}\right) + \frac{2b}{p_k}, \\ d_k &= c_k + \left(-\frac{n}{p_{k-1}} + d_{k-1}\right) \left(1 - \frac{2}{p_k}\right) + \frac{2}{p_k}, \\ \delta_k &= \delta_{k-1} \left(1 - \frac{2}{p_k}\right) + \frac{1}{p_k}, \quad \delta'_k = \delta'_{k-1} \left(1 - \frac{2}{p_k}\right), \\ \gamma_k &= \left(\gamma_{k-1} + \frac{n-1}{p_{k-1}}\right) \left(1 - \frac{2}{p_k}\right) - \frac{1}{p_k} - c_k. \end{aligned}$$

If p_k satisfies that $\tilde{p}_k < p_k < \infty$, where $\tilde{p}_k = 2 + \frac{2}{n-1} + \frac{2}{(n-1)^2} + \cdots + \frac{2}{(n-1)^{k-2}} + \frac{p_0}{(n-1)^{k-1}}$, then we can choose a, b , such that $c_k = 0$ and $d_k = -\frac{n-2}{2} + \frac{n}{p_k}$, $\gamma_k > 0$. Thus we have

$$\| |\nabla|^{d_k} \mathcal{D}_\omega^{\delta_k, \delta'_k} D_\sigma^{\gamma_k} u \|_{L_t^2 L_r^{p_k} L_\sigma^2} \leq \mathcal{C}_k \|\varphi\|_{L_x^2}. \quad (3.4)$$

3.2 A limiting argument

We first observe that \tilde{p}_k is decreasing and $\tilde{p}_k \rightarrow \frac{2(n-1)}{n-2} \equiv \tilde{p}$ as $k \rightarrow \infty$. Now we fix p with $\tilde{p} < p < p_0$ and let $k \rightarrow \infty$ to get the desired estimate (3.1).

Since $\tilde{p} < p < p_0$, $1 - \theta_k$ is bound away from 0 and 1, there exists $0 < \lambda < 1$ such that

$$\mathcal{C}_k \leq C_0^{1-\theta_k+(1-\theta_k)(1-\theta_{k-1})+\dots+\prod_{2 \leq j \leq k}(1-\theta_j)} C'_0 \mathcal{C}_1 \leq C_0^{\sum_{j \geq 1} \lambda^j} C'_0 \mathcal{C}_1 = C_0^{\frac{\lambda}{1-\lambda}} C'_0 \mathcal{C}_1.$$

If p is fixed near $\frac{2(n-1)}{n-2}$, then we can pick p_k such that $p_k = p$ for all $k \geq k_0$ for some large k_0 . Since $\delta_k = \delta_{k-1}(1 - 2/p) + 1/p$ for $k \geq k_0$, δ_k is increasing and bounded and thus $\lim_{k \rightarrow \infty} \delta_k = \frac{1}{2}$. Also $\delta'_k/\delta'_{k-1} = (1 - 2/p)$ implies $\delta'_k \rightarrow 0$ as $k \rightarrow \infty$. On the other hand, γ_k is decreasing and bounded below so that $\gamma_k \rightarrow \frac{n-2}{2} - \frac{n-1}{p}$. This limit goes to 0 as $p \rightarrow \frac{2(n-1)}{n-2}$.

By the usual density argument, for the proof of (3.1) we may assume $\hat{\varphi} \in C_0^\infty(\mathbb{R}^n \setminus \{0\})$. Since $|\nabla|^{d_k} \mathcal{D}_\omega^{\delta_k, \delta'_k} D_\sigma^{\gamma_k} u$ and $|\nabla|^{-\frac{n-2}{2} + \frac{n}{p}} \mathcal{D}_\omega^{\frac{1}{2}, 0} D_\sigma^{\frac{n-2}{2} - \frac{n-1}{p}} u$ are smooth and have the same compact Fourier support, it is easy to see that both are $O((|x|+1)^{-M}(|t|+1)^{-M})$ for any large M . Then it is obvious that

$$\lim_{k \rightarrow \infty} |\nabla|^{d_k} \mathcal{D}_\omega^{\delta_k, \delta'_k} D_\sigma^{\gamma_k} u = |\nabla|^{-\frac{n-2}{2} + \frac{n}{p}} \mathcal{D}_\omega^{\frac{1}{2}, 0} D_\sigma^{\frac{n-2}{2} - \frac{n-1}{p}} u.$$

Hence by taking limit

$$\lim_{k \rightarrow \infty} \| |\nabla|^{d_k} \mathcal{D}_\omega^{\delta_k, \delta'_k} D_\sigma^{\gamma_k} u \|_{L_t^2 L_r^{p_k} L_\sigma^2} = \| |\nabla|^{-\frac{n-2}{2} + \frac{n}{p}} \mathcal{D}_\omega^{\frac{1}{2}, 0} D_\sigma^{\frac{n-2}{2} - \frac{n-1}{p}} u \|_{L_t^2 L_r^p L_\sigma^2}.$$

Since \mathcal{C}_k is uniform on k , from (3.4) we get (3.1) provided that p is fixed near \tilde{p} . This completes the proof of Proposition 3.1.

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