A MAXIMAL INEQUALITY ASSOCIATED TO SCHRÖDINGER TYPE EQUATION

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ABSTRACT. In this note, we consider a maximal operator $\sup_{t\in\mathbb{R}}|u(x,t)|=\sup_{t\in\mathbb{R}}|e^{it\Omega(D)}f(x)|$, where u is the solution to the initial value problem $u_t=i\Omega(D)u,\ u(0)=f$ for a C^2 function Ω with some growth rate at infinity. We prove that the operator $\sup_{t\in\mathbb{R}}|u(x,t)|$ has a mapping property from a fractional Sobolev space $H^{\frac{1}{4}}$ with additional angular regularity in which the data lives to $L^2((1+|x|)^{-b}dx)\ (b>1)$. This mapping property implies the almost everywhere convergence of u(x,t) to f as $t\to 0$, if the data f has an angular regularity as well as $H^{\frac{1}{4}}$ regularity.

1. INTRODUCTION

We consider the following free Schrödinger type equation:

$$\frac{\partial}{\partial t}u(x,t) = i\Omega(D)u(x,t) \quad \text{in} \quad \mathbb{R}^{n+1} (n \ge 2), \quad u(x,0) = f(x),$$

where $\Omega(D)$ is a generalized differential operator defined by a C^2 function Ω and $D = (-\Delta)^{\frac{1}{2}}$. For smooth initial data f, the solution $u(x,t) = e^{it\Omega(D)}f$ can be written as

$$u(x,t) = \frac{1}{(2\pi)^n} \int_{\mathbb{R}^n} e^{i(x\cdot\xi + t\Omega(\xi))} \widehat{f}(\xi) \, d\xi, \quad f \in \mathcal{S}(\mathbb{R}^n),$$

where $\widehat{f}(\xi) = \int e^{-ix\cdot\xi} f(x) dx$. In this note, we assume that the initial data f has H^s regularity for some s > 0 as well as some regularity in the angular direction. For $\alpha, \beta \ge 0$, we define an initial data space $H_r^{\alpha} H_{\omega}^{\beta}$ by

$$H_{r}^{\alpha}H_{\omega}^{\beta} = \left\{ f: \|f\|_{H_{r}^{\alpha}H_{\omega}^{\beta}} := \|(1-\Delta)^{\frac{\alpha}{2}}f\|_{L_{r}^{2}H_{\omega}^{\beta}} < \infty \right\},$$

where $||g||_{L_r^2}^2 = \int_0^\infty |g(r)|^2 r^{n-1} dr$, $||g||_{L_r^2 H_\omega^\beta} = ||||(1 - \Delta_\omega)^{\frac{\beta}{2}} f(r\omega)||_{L_\omega^2}||_{L_r^2}$ (here, $(r, \omega) \in \mathbb{R}_+ \times S^{n-1}$ is the spherical coordinates), and Δ_ω is the Laplace-Beltrami operator on S^{n-1} . Since Δ_ω commutes with Δ , one can readily check that $||g||_{H_r^\alpha H_\omega^\beta} \sim ||(1 - \Delta_\omega)^{\frac{\beta}{2}} g||_{H^\alpha}$ (for instance, see [9]). Since not every function in $H_r^\alpha H_\omega^\beta$ has radial regularity higher than α , there is no embedding from or into a usual Sobolev

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space. In particular, it should be noted that $H_r^{\alpha} H_{\omega}^{\beta} \not\subseteq H^{\alpha+\gamma}$ $(0 < \gamma < \beta)$ and $H_r^{\alpha} H_{\omega}^{\beta} \not\supseteq H^{\alpha+\gamma}$ $(\gamma \ge \beta)$.

We also assume that $\Omega \in C^2(\mathbb{R}^n)$ is radially symmetric and satisfies

$$c_1|\rho|^{a-k} \le |\Omega^{(k)}(\rho)| \le c_2|\rho|^{a-k} \ (k=0,1,2), \quad \text{if} \quad |\rho| \ge N$$

for some $c_1, c_2, a > 0$ with $a \neq 1$ and a large N > 0. With the above assumptions, let us define a maximal function $u^*(x)$ by $u^*(x) = \sup_{t \in \mathbb{R}} |u(x,t)|$. We prove

Theorem 1.1. For any $\varepsilon > 0$ and b > 1, if $f \in H_r^{\frac{1}{4}} H_{\omega}^{\frac{n-1}{2} - \frac{1}{4} + \varepsilon}$, then there exists a constant C, depending only on $a, c_1, c_2, N, n, \varepsilon, b$, such that

$$||u^*||_{L^2((1+|x|)^{-b}dx)} \le C||f||_{H_r^{\frac{1}{4}}H_\omega^{\frac{n-1}{2}-\frac{1}{4}+\varepsilon}}.$$

Now let us define a linear operator T and a maximal operator T^* for a fixed s > 0 by

$$Tf(x,t) = w(|x|) \int e^{i(x\cdot\xi + t\Omega(\xi))} \hat{f}(\xi) \, \frac{d\xi}{(1+|\xi|^2)^{\frac{s}{2}}}$$

where $w(r) = (1+r)^{-\frac{b}{2}}, b > 0$ and

$$T^*f(x) = \sup_{t \in \mathbb{R}} |Tf(x,t)|.$$

Then Theorem 1.1 follows immediately from

Theorem 1.2. For any $\varepsilon > 0$ and b > 1, if $f \in L^2_r H^{\frac{n-1}{2}-s+\varepsilon}_{\omega}$ for some $s \in [\frac{1}{4}, \frac{1}{2})$, there exists a constant C, depending only on $a, c_1, c_2, N, n, s, \varepsilon, b$, such that

$$||T^*f||_{L^2} \le C||f||_{L^2_r H^{\frac{n-1}{2}-s+\varepsilon}_{\omega}}.$$

The maximal function u^* and operator T^* have been studied extensively by many authors ([1, 2, 3, 4, 5, 7, 8, 10, 11, 12, 13, 14, 18, 19, 21]). P. Sjölin [14] and L. Vega [19] showed that for some ball B_R of radius R

(1.1)
$$||u^*||_{L^2(B_R)} \le C||f||_{H^s},$$

only if $s \ge \frac{1}{4}$. Up to now, it is known that (1.1) is true if n = 1 ([5, 8]) or the initial data is radial ([4, 12]), or $s > \frac{1}{2}$ and $n \ge 2$ ([11, 19]). Recently, T. Tao [18] obtained (1.1) for $s > \frac{2}{5}$ and n = 2. However, the sufficiency remains open widely.

On the other hand, Theorem 1.1 shows that it is true for $s = \frac{1}{4}$ if we assume the additional angular regularity. If the initial data is a finite linear combination of

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radial functions and spherical harmonics such that $f = \sum_{k \leq L} f_k Y_k$, it was proved by the first and third authors in [4] that $||u^*||_{L^{\frac{4n}{2n-1}}} \leq C_L ||f||_{H^{\frac{1}{4}}}$, where

$$C_L \le CL^{\frac{1}{2} + \varepsilon} (n + 2L)^{\frac{n+2L}{2}} \max_{1 \le k \le L} \frac{\|Y_k\|_{L^{\frac{4n}{2n-1}}}}{\|Y_k\|_{L^2}} \quad (0 < \varepsilon \ll 1)$$

The factor $(n + 2L)^{\frac{n+2L}{2}}$ follows from the asymptotic behavior of Bessel function $(J_{\nu}(t) \sim b_{+}t^{-\frac{1}{2}}e^{it} + b_{-}t^{-\frac{1}{2}}e^{-it} + O((n+2\nu)^{\frac{n+2\nu}{2}})t^{-\frac{3}{2}}$ for t > 1). The tail $t^{-\frac{3}{2}}$ seems inevitable to obtain the non-weighted global $L^{\frac{4n}{2n-1}}$ ($\frac{4n}{2n-1} > 2$) estimate for which a big cost of C_{L} is paid. In this connection, Theorem 1.1 improves significantly the dependency on the order of spherical harmonic up to $L^{3/4+\varepsilon}$ (see (2.2) below). This improvement occurs from an estimate for the tail of Bessel function Ct^{-1} for $t > 2\nu$, which enables us to use the L^{2} method. The weighted L^{2} estimate as in Theorem 1.1 is absolutely necessary for a global estimate in view of the negative result that the non-weighted global L^{2} estimate [11] and any local estimate in $L^{p}(p > 2)$ [22] are impossible for the data $f \in H^{\frac{1}{4}}$.

In case that $\Omega(D) = -\Delta$, recently G. Gigante and F. Soria [6] showed a local L^2 estimate, independently of our work, that $||u^*||_{L^2(B_R)} \leq CL^{\frac{1}{2}+\varepsilon}||f||_{H^{\frac{1}{4}}}$. They used a finer asymptotic behavior of Bessel function $J_{\nu}(t)$ for $\nu + \nu^{\frac{1}{3}} \leq t \leq 2\nu$ but their method seems not to be applied directly to the general phase Ω because the power of L may depend on Ω in their argument.

From the assumption on Ω , we treat Ω not only of the form $|\xi|^a$ but also $\sum_{i=1}^l m_i |\xi|^{a_i}$ for any number $a_l > a_{l-1} > \cdots > a_1 > 0$, $a_l \neq 1$ and $m_i \in \mathbb{R}$. For the more general phase Ω , we refer the readers to [3] in which a weighted L^2 estimate is discussed for the phase Ω with $\nabla \Omega$ having zeros or singularities. For another use of angular regularity, one can refer to [9] in which the endpoint Strichartz estimates of 3-d wave and Klein-Gordon equations are considered.

If not specified, throughout this paper, C denotes a generic constant that depends on $a, c_1, c_2, N, n, s, b, \varepsilon$. We use the notation $A \leq B$ and $A \sim B$ to denote $|A| \leq CB$ and $C^{-1}B \leq |A| \leq CB$ respectively.

2. Proof of Theorem 1.2

We begin with reviewing some properties of the spherical harmonic expansion. If $f(r\omega) = g(r)Y_k(\omega)$ for a radial function g and a spherical harmonic Y_k of order k, then we have

$$f(\rho\theta) = G(\rho)Y_k(\theta), \quad ||g||_{L^2_r} = ||G||_{L^2_r},$$

where

$$G(\rho) = c_{n,k} \int_0^\infty g(r) r^{n-1} (r\rho)^{-\frac{n-2}{2}} J_\nu(r\rho) \, dr, \ |c_{n,k}| \le C, \ \nu = \frac{2k+n-2}{2}.$$

For the representation of G, see e.g. [16] or [22]. Since $-\Delta_{\omega}Y_k = k(k+n-2)Y_k$, we also have $||f||_{L^2_r H^{\beta}_{\omega}} \sim (1+k^2)^{\frac{\beta}{2}} ||g||_{L^2_r} ||Y_k||_{L^2_{\omega}}$. Furthermore, if $h \in L^2_r H^{\beta}_{\omega}$, then there exist radial functions $\{h^l_k\}$ and spherical harmonics $\{Y^l_k\}$ such that

$$h(r\omega) = \sum_{k \ge 0} \sum_{1 \le l \le d(k)} h_k^l(r) Y_k^l(\omega) \quad \text{in} \quad L_r^2 H_\omega^\beta,$$

where d(k) is the dimension of the space of spherical harmonics of degree k, and

(2.1)
$$||h||_{L^2_r H^\beta_\omega}^2 \sim \sum_{k \ge 0} \sum_{1 \le l \le d(k)} (1+k^2)^\beta ||h^l_k||_{L^2_r}^2 ||Y^l_k||_{L^2_\omega}^2.$$

Thus for the proof of theorem, we have only to consider the case that $f(r\omega) = g(r)Y_k(\omega)$ and to show that for large k

(2.2)
$$||T^*f||_{L^2} \lesssim k^{\frac{1}{2}-s} ||g||_{L^2_r} ||Y_k||_{L^2_{\omega}},$$

since for the function $h(r\omega) = \sum_{k\geq 0} \sum_{1\leq l\leq d(k)} h_k^l(r) Y_k^l(\omega)$ in $L_r^2 H_{\omega}^{\beta}$, we have from (2.2)

$$\begin{split} \|T^*h\|_{L^2} &\lesssim \sum_k \sum_{1 \le l \le d(k)} k^{\frac{1}{2}-s} \|h_k^l\|_{L^2_r} \|Y_k^l\|_{L^2_\omega} \\ &\lesssim \sum_k k^{\frac{1}{2}-s} d(k)^{\frac{1}{2}} \left(\sum_{1 \le l \le d(k)} \|h_k^l\|_{L^2_r}^2 \|Y_k^l\|_{L^2_\omega}^2 \right)^{\frac{1}{2}} \\ &\lesssim \sum_k k^{\frac{n-1}{2}-s} \left(\sum_{1 \le l \le d(k)} \|h_k^l\|_{L^2_r}^2 \|Y_k^l\|_{L^2_\omega}^2 \right)^{\frac{1}{2}} \\ &\lesssim \left(\sum_k \sum_{1 \le l \le d(k)} k^{n-1-2s+\varepsilon} \|h_k^l\|_{L^2_r}^2 \|Y_k^l\|_{L^2_\omega}^2 \right)^{\frac{1}{2}}, \end{split}$$

where we used the estimate $d(k) = \frac{n+2k-2}{k} \binom{n+k-3}{k-1} \lesssim k^{n-2}$ for the third inequality (see [16]).

Now if $\widehat{f}(\rho\omega) = G(\rho)Y_k(\omega)$, from the definition of T, it follows that

$$Tf(r\omega,t) = w(r) \int_{S^{n-1}} \int_0^\infty e^{i(r\omega \cdot \rho \theta + t\Omega(\rho))} G(\rho) Y_k(\theta) \rho^{n-1} \frac{d\rho}{(1+\rho^2)^{\frac{s}{2}}} d\theta$$

= $c_{n,k} w(r) \int_0^\infty e^{it\Omega(\rho)} (r\rho)^{-\frac{n-2}{2}} J_\nu(r\rho) \rho^{n-1} G(\rho) \frac{d\rho}{(1+\rho^2)^{\frac{s}{2}}} Y_k(-\omega).$

Let us define an operator S by

$$SG(r,t) = c_{n,k} r^{\frac{n-1}{2}} w(r) \int_0^\infty e^{it\Omega(\rho)} (r\rho)^{-\frac{n-2}{2}} J_\nu(r\rho) \rho^{\frac{n-1}{2}} G(\rho) \frac{d\rho}{(1+\rho^2)^{\frac{s}{2}}}.$$

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Let us denote by $||F||_{L^pL^q}$ the mixed norm $\|(\|F(r,t)\|_{L^q(dt)})\|_{L^p(dr)}$. Here we use the notation $||F||_{L^p(dr)}^p$ for $\int |F(r)|^p dr$ to avoid the confusion with $||F||_{L^p}$. To prove (2.2) it suffices to show that

(2.3)
$$\|S\tilde{G}\|_{L^{2}L^{\infty}} \lesssim k^{\frac{1}{2}-s} \|\tilde{G}\|_{L^{2}(dr)},$$

where $\tilde{G}(\rho) = \rho^{\frac{n-1}{2}} G(\rho)$. Now we define the dual operator S^d of S by

$$S^{d}F(\rho) = \frac{c_{n,k}}{(1+\rho^{2})^{\frac{s}{2}}} \int_{\mathbb{R}} \int_{0}^{\infty} e^{-it\Omega(\rho)} (r\rho)^{\frac{1}{2}} J_{\nu}(r\rho) w(r) F(r,t) \, dr dt$$

for $F \in C_0^{\infty}(\mathbb{R}_+ \times \mathbb{R})$. Then, by duality (2.3) follows from

(2.4)
$$||S^{d}F||_{L^{2}(dr)} \leq Ck^{\frac{1}{2}-s}||F||_{L^{2}L^{1}}.$$

Choose smooth cut-off functions ϕ_0 , ϕ_1 and ϕ_3 so that $\phi_0 = 1$ on $\{|s| < \frac{1}{4}\}$, $\phi_0 = 0$ on $\{|s| > \frac{1}{2}\}$, $\phi_1 = 1$ on $\{|s| \sim 1\}$, $\phi_1 = 0$ otherwise, $\phi_2 = 0$ on $\{|s| < 2\}$, $\phi_2 = 1$ on $\{|s| > 3\}$, and $\phi_0 + \phi_1 + \phi_2 = 1$. Then we decompose S^d as

$$S^d F(\rho) = S_0 F + S_1 F + S_2 F,$$

where for i = 0, 1, 2,

$$S_i F(\rho) = \frac{c_{n,k}}{(1+\rho^2)^{\frac{s}{2}}} \int_{\mathbb{R}} \int_0^\infty e^{-it\Omega(\rho)} (r\rho)^{\frac{1}{2}} J_\nu(r\rho) \phi_i\left(\frac{r\rho}{\nu}\right) w(r) F(r,t) \, dr dt.$$

Now we need to show each S_i satisfies (2.4) in the place of S^d . Each estimate is to be shown using the following asymptotic behavior of Bessel functions:

(2.5)
$$|J_{\nu}(t)| \le C \exp(-C\nu), \text{ if } t \le \frac{\nu}{2},$$

(2.6)
$$\frac{1}{r} \int_0^r |J_{\nu}(t)|^2 t \, dt \le C \quad \text{for all } r > 0,$$

(2.7)
$$J_{\nu}(t)\phi_{2}(\frac{t}{\nu}) = t^{-\frac{1}{2}}(b_{+}e^{it} + b_{-}e^{-it})\phi_{2}(\frac{t}{\nu}) + \Phi_{\nu}(t)\phi_{2}(\frac{t}{\nu})$$

where $|\Phi_{\nu}(t)| \leq \frac{C}{t}$, $|b_{\pm}| \leq C$ and the constant C is independent of ν . For the proof of (2.5), see [17]. The mean value estimate (2.6) can be found in Section 4.10 of [20]. Invoking the Schläfli's integral representation (see p.176 in [23]):

$$J_{\nu}(t) = \frac{1}{2\pi} \int_{0}^{2\pi} e^{i(t\sin\theta - \nu\theta)} d\theta - \frac{\sin(\nu\pi)}{\pi} \int_{0}^{\infty} e^{-\nu\tau - t\sinh\tau} d\tau,$$

the last two asymptotic behavior (2.7) follow from the easy estimate

$$\left|\frac{\sin(\nu\pi)}{\pi}\int_0^\infty e^{-\nu\tau-t\sinh\tau}\,d\tau\right| \le \frac{C}{\nu+t}$$

and the method of stationary phase such that

$$\frac{1}{2\pi} \int_0^{2\pi} e^{i(t\sin\theta - \nu\theta)} d\theta = (b_+ e^{it} + b_- e^{-it})t^{-\frac{1}{2}} + O(t^{-\frac{3}{2}}) \quad \text{for} \quad t > 2\nu.$$

Using (2.5), we now see

$$\begin{split} |S_0 F(\rho)| &\lesssim \nu^{\frac{1}{2}} e^{-C\nu} (1+\rho^2)^{-\frac{s}{2}} \int_0^{\frac{\nu}{\rho}} w(r) ||F(r,\cdot)||_{L^1} dr \\ &= \nu^{\frac{1}{2}} e^{-C\nu} (1+\rho^2)^{-\frac{s}{2}} \left(\int_0^{\min(\frac{\nu}{\rho},2)} ||F(r,\cdot)||_{L^1} dr \\ &+ \int_0^{\frac{\nu}{\rho}} \chi_{[2,\infty)}(r) w(r) ||F(r,\cdot)||_{L^1} dr \right) \\ &\lesssim \nu^{\frac{1}{2}} e^{-C\nu} (1+\rho^2)^{-\frac{s}{2}} \left((\min(\frac{\nu}{\rho},2))^{\frac{1}{2}} + \chi_{[0,\frac{\nu}{2}](\rho)} \right) ||F||_{L^2 L^1}. \end{split}$$

Thus we have

 $||S_0F||_{L^2(dr)}$

(2.8)
$$\lesssim \nu^{\frac{1}{2}} e^{-C\nu} \left(\int_0^\infty (1+\rho^2)^{-s} (\min(\frac{\nu}{\rho}, 2) + \chi_{[0,\frac{\nu}{2}]}(\rho)) \, d\rho \right)^{\frac{1}{2}} ||F||_{L^2 L^1} \\ \lesssim \nu^{1-s} e^{-C\nu} ||F||_{L^2 L^1}.$$

For S_1 , we have

$$|S_1 F(\rho)| \lesssim (1+\rho^2)^{-\frac{s}{2}} \left(\int_0^\infty J_\nu^2(r\rho) r\rho \phi_1^2\left(\frac{r\rho}{\nu}\right) w(r)^2 dr \right)^{\frac{1}{2}} ||F||_{L^2 L^1}$$
$$\lesssim (1+\rho^2)^{-\frac{s}{2}} \left(\int_0^2 + \int_2^\infty \right)^{\frac{1}{2}} ||F||_{L^2 L^1}.$$

Using the change of variable $r \mapsto r/\rho$, the first part in the middle parenthesis is bounded by $\chi_{[\frac{\nu}{4},\infty)}(\rho)\frac{1}{\rho}\int_{0}^{2\rho}J_{\nu}^{2}(r)r\phi_{1}^{2}(r/\nu) dr$. By (2.6), it follows that

$$\int_0^2 \lesssim \nu \rho^{-1} \chi_{\left[\frac{\nu}{4},\infty\right)}(\rho).$$

For the second part, we also use the change of variable $r\mapsto r/\rho$ and then by (2.6) have that

$$\int_{2}^{\infty} \lesssim \rho^{b-1} \int_{\max(2\rho, \frac{\nu}{2})}^{3\nu} J_{\nu}^{2}(r) r^{1-b} dr \lesssim \nu \rho^{b-1}(\max(2\rho, \frac{\nu}{2}))^{-b}.$$

We thus obtain

$$||S_1F||_{L^2(dr)}$$

(2.9)
$$\lesssim \left(\int_0^\infty (1+\rho^2)^{-s} (\nu\rho^{-1}\chi_{[\frac{\nu}{4},\infty)}(\rho) + \nu\rho^{b-1}(\max(2\rho,\frac{\nu}{2}))^{-b}) \, d\rho \right)^{\frac{1}{2}} ||F||_{L^2L^1} \\ \lesssim \nu^{\frac{1}{2}-s} ||F||_{L^2L^1}.$$

Now we estimate S_2F . Let us set $S_2F = S_+F + S_-F + S_3F$, where

$$S_{\pm}F(\rho) = \frac{c_{n,k}b_{\pm}}{(1+\rho^2)^{\frac{s}{2}}} \int_{\mathbb{R}} \int_0^{\infty} e^{i(\pm r\rho - t\Omega(\rho))} \phi_2(\frac{r\rho}{\nu}) w(r) F(r,t) \, dr dt$$

$$S_3F(\rho) = \frac{c_{n,k}}{(1+\rho^2)^{\frac{s}{2}}} \int_{\mathbb{R}} \int_0^{\infty} e^{-it\Omega(\rho)} (r\rho)^{\frac{1}{2}} \Phi_{\nu}(r\rho) \phi_2\left(\frac{r\rho}{\nu}\right) w(r) F(r,t) \, dr dt.$$

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For the estimate $S_{\pm}F$, it suffices to consider $S_{+}F$. We decompose it into two parts as follows:

$$S_+F(\rho) = I + II$$

where

$$I = \frac{c_{n,k}b_{+}}{(1+\rho^{2})^{\frac{s}{2}}} \int_{\mathbb{R}} \int_{0}^{\infty} e^{i(r\rho-t\Omega(\rho))} w(r)F(r,t) \, dr dt,$$

$$II = \frac{c_{n,k}b_{+}}{(1+\rho^{2})^{\frac{s}{2}}} \int_{\mathbb{R}} \int_{0}^{\infty} e^{i(r\rho-t\Omega(\rho))} (\phi_{2}(\frac{r\rho}{\nu}) - 1)w(r)F(r,t) \, dr dt.$$

For II, we have

$$|II(\rho)| \lesssim (1+\rho^2)^{-\frac{s}{2}} \int_0^{\frac{3\nu}{\rho}} w(r) ||F(r,\cdot)||_{L^1} dr$$
$$\lesssim (1+\rho^2)^{-\frac{s}{2}} \left(\int_0^{\frac{3\nu}{\rho}} w(r)^2 dr \right)^{\frac{1}{2}} ||F||_{L^2L^1}$$

and hence by the similar estimate to (2.8) for S_0F

(2.10)
$$||II||_{L^2(dr)} \lesssim \nu^{\frac{1}{2}-s} ||F||_{L^2L^1}$$

Now we estimate I. Since F is in $C_0^{\infty}(\mathbb{R}_+ \times \mathbb{R})$, obviously we may assume

$$I = \frac{c_{n,k}b_+}{(1+\rho^2)^{\frac{s}{2}}} \int_{\mathbb{R}^2} e^{i(r\rho - t\Omega(\rho))} w(|r|) F(r,t) \, dr dt.$$

Squaring and integrating I over $\{|\rho| \leq N\}$, where N is the number in the condition of Ω , we have

(2.11)
$$\int_{|\rho| < N} |I|^2 \, d\rho \le C ||F||_{L^2 L^1}^2$$

Now it is easy to see

$$\int_{|\rho| > N} |I|^2 d\rho$$

$$\leq C \iiint |K(r - r', t - t')w(|r|)|F(r, t)|w(|r'|)|F(r', t')| dr dr' dt dt',$$

where

$$K(r,t) = \int_{|\rho| > N} e^{i(r\rho - t\Omega(\rho))} \frac{d\rho}{|\rho|^{2s}}.$$

For the kernel estimate, we introduce a lemma which shows uniform bound of kernel K on t.

Lemma 2.1 (see Lemma 2.3 in [4]). For any real number $A, B(A \neq 0)$ and $s \in [\frac{1}{2}, 1)$, there exists a constant C independent of A and B such that

$$\left| \int_{|\rho| > N} e^{i(A\Omega(\rho) + B\rho)} \frac{d\rho}{|\rho|^s} \right| \le C|B|^{-(1-s)}.$$

Applying Lemma 2.1 with 2s $(\frac{1}{4} \leq s < \frac{1}{2})$ and B = r - r', from fractional integration and Hölder inequality it follows

$$(2.12) \qquad \int_{|\rho|>N} |I|^2 d\rho \\ \lesssim \iint |r-r'|^{-(1-2s)} w(|r|) ||F(r,\cdot)||_{L^1} w(|r'|) ||F(r',\cdot)|| \, dr dr \\ \lesssim ||\mathcal{I}_{2s}(w||F||_{L^1})||_{L^p} ||w||F||_{L^1} ||_{L^{p'}} \quad \left(\frac{1}{p} = \frac{1}{p'} - 2s\right) \\ \lesssim ||wF||_{L^{\frac{2}{1+2s}}L^1}^2 \lesssim ||w||_{L^{\frac{1}{s}}}^2 ||F||_{L^{2}L^1}^2 \quad \left(\frac{b}{2} \cdot \frac{1}{s} > 1\right) \\ \lesssim ||F||_{L^{2}L^1}^2,$$

where \mathcal{I}_{2s} is the Riesz potential of order 2s.

Finally, we estimate S_3F . From the uniform bound of Φ_{ν} on ν , for small $\varepsilon > 0$, we have

$$\begin{split} &|S_{3}F(\rho)|\\ \lesssim \frac{1}{(1+\rho^{2})^{\frac{s}{2}}} \int (r\rho)^{-\frac{1}{2}} \phi_{2}\left(\frac{r\rho}{\nu}\right) w(r) ||F(r,\cdot)||_{L^{1}} dr\\ \lesssim \rho^{-s-\frac{1}{2}} \chi_{[\nu,\infty)}(\rho) \int_{\frac{2\nu}{\rho}}^{2} r^{-\frac{1}{2}} ||F(r,\cdot)||_{L^{1}} dr + \rho^{-s-\frac{1}{2}} \int_{\max(2,\frac{2\nu}{\rho})} r^{-\frac{1}{2}-\frac{b}{2}} ||F(r,\cdot)||_{L^{1}} dr\\ \lesssim \nu^{-\delta} \rho^{-s-\frac{1}{2}+\delta} \chi_{[\nu,\infty)}(\rho) \int_{\frac{2\nu}{\rho}}^{2} r^{-\frac{1}{2}+\delta} ||F(r,\cdot)||_{L^{1}} dr + \rho^{-s-\frac{1}{2}} \left(\max(2,\frac{2\nu}{\rho})\right)^{-\frac{b}{2}} ||F||_{L^{2}L^{1}}\\ \lesssim \left(\nu^{-\delta} \rho^{-s-\frac{1}{2}+\delta} \chi_{[\nu,\infty)}(\rho) + \rho^{-s-\frac{1}{2}} \left(\max(2,\frac{2\nu}{\rho})\right)^{-\frac{b}{2}}\right) ||F||_{L^{2}L^{1}}. \end{split}$$

Choosing δ as $\frac{1}{8}$, we obtain

(2.13)
$$||S_3F||_{L^2(dr)} \lesssim \nu^{-s} ||F||_{L^2L^1}.$$

Combining all the estimates (2.8) to (2.13) and recalling $\nu = \frac{2k+n-2}{2}$, we get (2.4) and hence Theorem 1.2.

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References

- [1] J. Bourgain, A remark on Schrödinger operators, Israel J. Math. 77 (1992), 1-16.
- [2] L. Carleson, Some analytical problems related to statistical mechanics, Euclidean Harmonic Analysis, Lecture Notes in Math. 779 (1979), 5-45.
- [3] Y. Cho and Y. Shim, Weighted L² estimates for maximal operators associated to dispersive equation, Illinois J. Math. 48 (2004), 1081-1092.

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- [4] Y. Cho and Y. Shim, Global estimates of maximal operators generated by dispersive equations, Hokkaido Univ. Preprint Series in Mathematics #704, (http://eprints.math.sci.hokudai.ac.jp/archive/00000883/).
- [5] B. E. J. Dahlberg and C. E. Kenig, A note on almost everywhere behavior of solutions to the Schrödinger equation, Harmonic Analysis, Lecture Notes in Math. 908 (1982), 205-209.
- [6] G. Gigante and F. Soria, On the boundedness in H^{1/4} of the maximal square function associated with the Schrödinger equation, in preprint.
- [7] H. P. Heinig and S. Wang, Maximal function estimates of solutions to general dispersive partial differential equations, Trans. Amer. Math. Soc. (1) 351 (1999), 79-108.
- [8] C. E. Kenig and A. Ruiz, A strong type (2,2) estimate for a maximal operator associated to the Schrödinger equation, Trans. Amer. Math. Soc. 280 (1983), 239-246.
- [9] S. Machihara, M. Nakamura, K. Nakanishi and T. Ozawa, Endpoint Strichartz estimates and global solutions for the noninear Dirac equation, J. Func. Anal. 219 (2005), 1-20.
- [10] A. Moyua, A. Vargas and L. Vega, Restriction theorems and Maximal operators related to oscillatory integrals in ℝ³, Duke Math. J. (3) 96 (1999), 547-574.
- [11] P. Sjölin, Global maximal estimates for solutions to the Schrödinger equation, Studia. Math.
 (2) 110 (1994), 105-114.
- P. Sjölin, Radial functions and maximal estimates for solutions to the Schrödinger equation, J. Austral. Math. Soc. (Series A) 59 (1995), 134-142.
- [13] P. Sjölin, L^p Maximal estimates for solutions to the Schrödinger equation, Math. Scand. 81 (1997), 35-68.
- [14] P. Sjölin, A Counter-example Concerning Maximal Estimates for Solutions to Equations of Schrödinger Type, Indiana Univ. Math. J. (2) 47 (1998), 593-599.
- [15] E. M. Stein, Harmonic Analysis:Real-Variable Methods, Orthogonality, and Oscillatory Integrals, Princeton Univ. Press, Princeton, N.J., (1993).
- [16] E. M. Stein and G. Weiss, Introduction to Fourier Analysis on Euclidean Spaces, Princeton Univ. Press, (1971).
- [17] K. Stempak, A weighted uniform L^p- estimate of Bessel functions: a note on a paper of Guo, Proc. Amer. Math. Soc. 128 (2000), 2943-2945.
- [18] T. Tao, A sharp bilinear restriction estimate for paraboloids, Geom. Funct. Anal. 13 (2003), 1359-1384.
- [19] L. Vega, Schrödinger equations: pointwise convergence to the initial data, Proc. Amer. Math. Soc. 102 (1988),874-878.
- [20] B. G. Walther, Homogeneous estimates for oscillatory integrals, Acta Math. Univ. Comenianae, 9 (2000), 151-171.
- [21] B. G. Walther, Higher integrability for maximal oscillatory Fourier integrals, Annales Academiæ Scientiarum Fennicæ Mathematica 26 (2001), 189-204.
- [22] S. Wang, On the maximal operator associated with the free Schrödinger equation, Studia Math. 122 (1997), 167-182.
- [23] G. Watson, A treatise on the theory of Bessel functions, Reprint of the second (1944) edition. Cambridge Mathematical Library. Cambridge University Press, Cambridge, 1995.

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